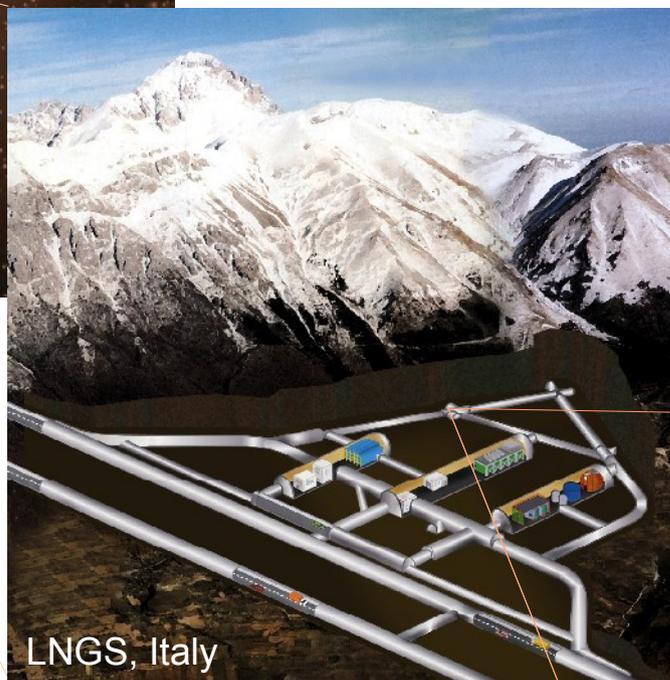
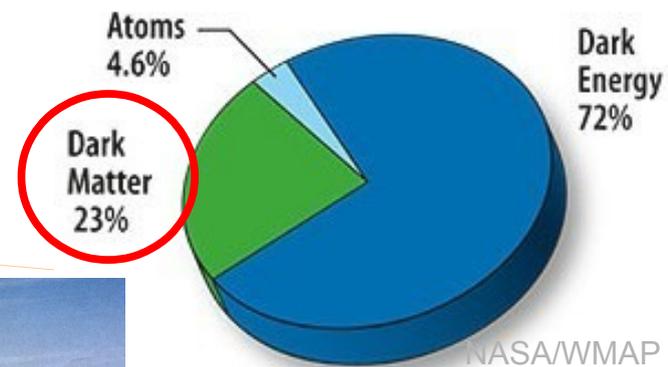


# Searches for Dark Matter



LNGS, Italy



XENON100 TPC

Lecture for *Astroteilchenschule*  
Obertrubach-Bärnfels  
October 11-13, 2011

Uwe Oberlack



JOHANNES GUTENBERG  
UNIVERSITÄT MAINZ

**Part 2:  
Fundamentals of Direct Detection**

# Outline of Lectures at Astroteilchenschule

## Part 1

- Evidence for Dark Matter
- WIMP Dark Matter
- Accelerator Searches
- Indirect Searches

# Outline of Lectures at Astroteilchenschule

## Part 2

- Direct Detection Technique
  - Kinematics
  - Energy Spectrum
  - Astro, Nuclear, Particle Physics Inputs
- Experimental backgrounds
- Detector techniques:
  - Noble liquids
  - Cryogenic germanium
  - Cryogenic scintillating crystals
  - Superheated liquids

# Outline of Lectures at Astroteilchenschule

## Part 3

- Signals (?)
  - DAMA / LIBRA annual modulation
  - CoGeNT
  - CRESST-II
- and Limits
  - CDMS-II
  - EDELWEISS-II
  - COUPP
  - XENON100
- Future

# Kinematics of DM Direct Detection

- Elastic scattering of WIMP  $\chi$  with a nucleus of mass  $m_n$  (given by atomic mass A).
- Measure energy  $E_R$  of recoiling nucleus:

$$E_R = \frac{\mu^2 v^2}{m_N} (1 - \cos \theta_{CM}) = 4 \frac{m_\chi m_N}{(m_\chi + m_N)^2} E_\chi \cos^2 \theta_N$$

$$E_{R,max} = 4 \frac{m_\chi m_N}{(m_\chi + m_N)^2} E_\chi = r E_\chi = 2 \frac{\mu^2}{m_\chi} v^2$$

$$v_{min} = \left( \frac{m_\chi E_{th}}{2\mu^2} \right)^{1/2} \quad \text{min. velocity to produce a signal above threshold}$$

- Useful abbreviations:

▶ **reduced mass**:  $\mu = (m_\chi m_N)/(m_\chi + m_N)$

$m_\chi \ll m_N$ :  $\mu \rightarrow m_\chi$        $m_\chi = m_N$ :  $\mu = m_\chi / 2$

$m_\chi \gg m_N$ :  $\mu \rightarrow m_N$

▶ **kinematic factor**  $r = 4 (m_\chi m_N)/(m_\chi + m_N)^2 = 4 \mu^2 / (m_\chi m_N)$

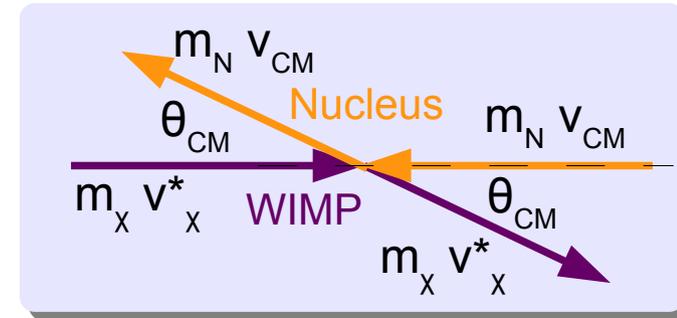
- Example:

▶  $m_\chi \sim 10 \text{ GeV}/c^2$ ,  $m_n = 135 \text{ GeV}/c^2$

▶  $v \sim 230 \text{ km/s}$ ,  $\beta = 7.7e-4$

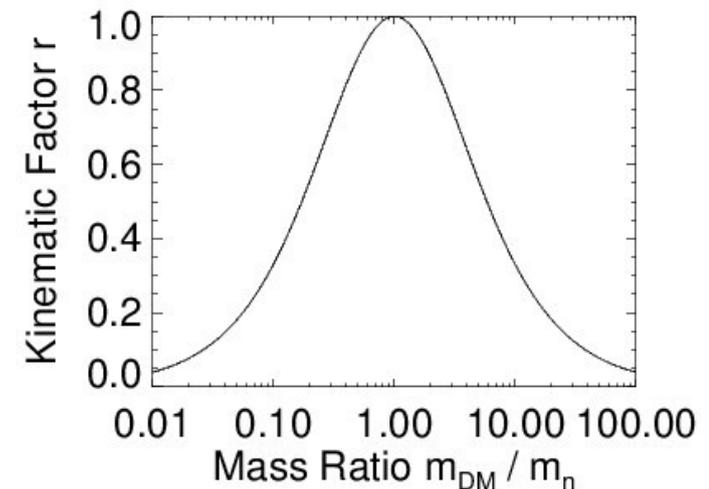
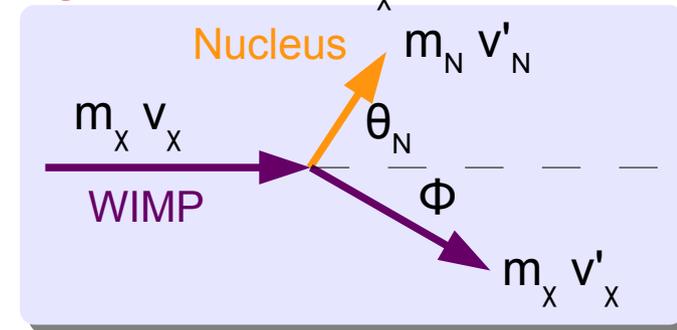
▶  $E_\chi = 29.4 \text{ keV}$ ,  $E_{R,max} = 7.5 \text{ keV}$

**CM system**



**Laboratory system**

short:  $v = v_\chi$



# WIMP Dark Matter Direct Detection

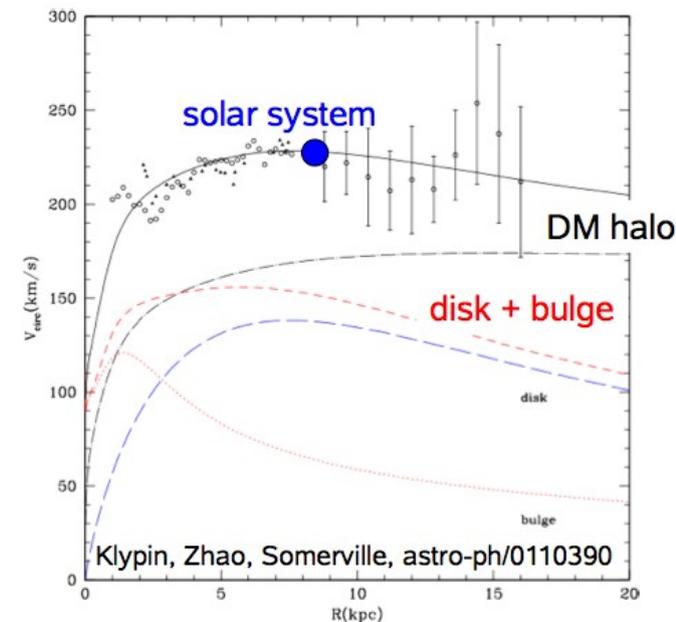
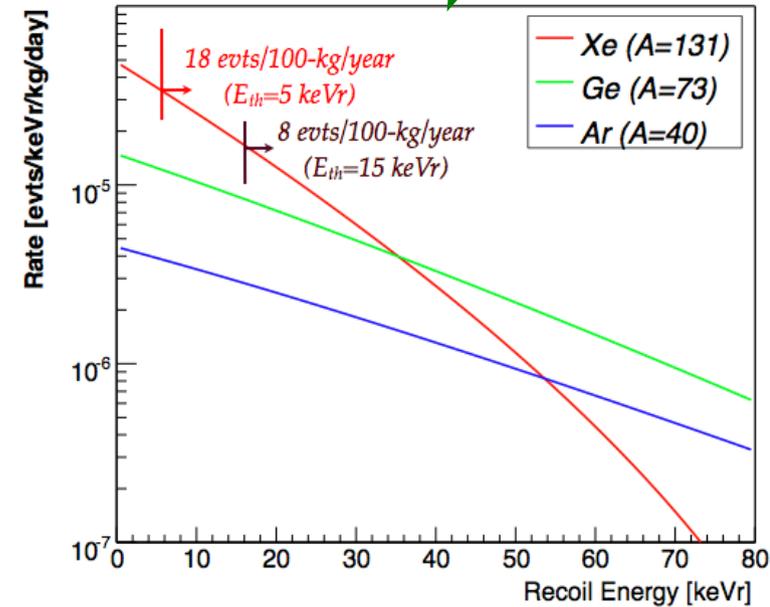
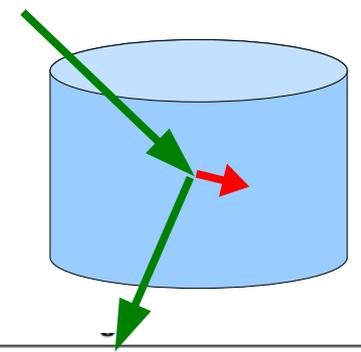
- Scattering of WIMPs  $\chi$  off of nuclei A.
  - ▶ elastic or inelastic?
  - ▶ spin-independent ( $\sim A^2$ ) or spin-dependent?
- Differential rate per unit detector mass:

$$\frac{dR}{dE} = \frac{\rho_\chi \sigma_s}{2 m_\chi \mu^2} |F(E)|^2 \int_{v_{min}}^{v_{esc}} f(\mathbf{v}, t) \frac{d^3 v}{v}$$

$$f(\mathbf{v}, t) \propto \exp\left(-\frac{(\mathbf{v} + \mathbf{v}_E(t))^2}{2\sigma_v^2}\right)$$

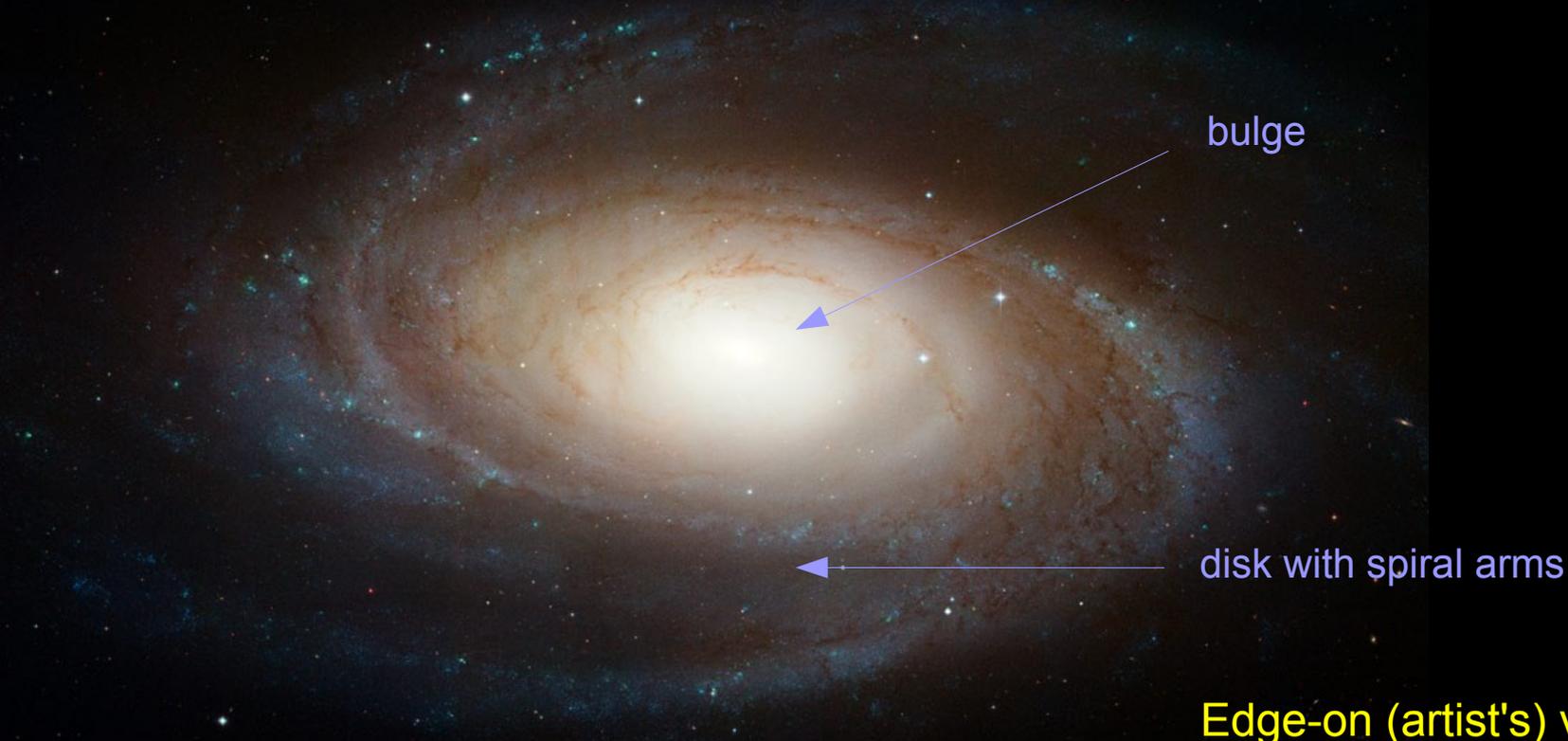
- ▶  $m_\chi \sim 10 - 10^4 \text{ GeV}/c^2$ ,  $\mu = (m_\chi m_N)/(m_\chi + m_N)$
- ▶  $v_\chi \sim 230 \text{ km/s}$
- ▶ “Standard” spherical halo: Featureless recoil spectrum  $\langle E \rangle \sim O(10 \text{ keV})$
- ▶  $\rho_\chi/m_\chi$ : local number density of WIMPs
- ▶  $\rho_\chi \sim 0.3 \text{ GeV}/c^2/\text{cm}^3$ ,  $\rho_\chi/m_\chi \lesssim 10 / L$
- ▶  $\sigma_s$  scattering cross section per nucleus.

Typical rate  $< 10^{-2}$  events / kg / day



# Astrophysics Input

# Structure of our Milky Way Galaxy

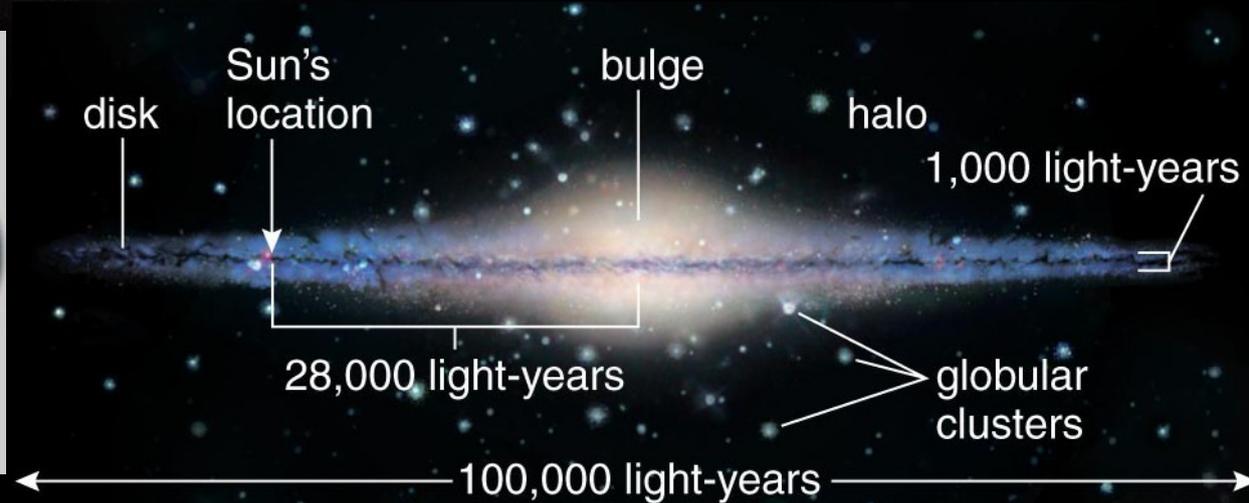


## Edge-on (artist's) view of the Milky Way

M81 Spiral Galaxy. NASA/HST



Milky Way (2MASS)



# Stellar Orbits in the Galaxy

*Halo stars travel high above and far below the disk on orbits with random orientations.*

*Bulge stars also have orbits with random orientations.*

*Disk stars orbit in circles with the same orientation, except for a little up-and-down motion.*

**Solar System:**

- ~15 pc above the galactic plane within the disk
- ~8.0 kpc from the Galactic Center

Pearson Education

$z=0.0$

Via Lactea 2 (2008)  
<http://www.ucolick.org/~diemand/vl>

# WIMP Dark Matter Direct Detection



Observer

80 kpc



# Halo Models

Considering the standard halo model, an isothermal sphere, the WIMP velocity distribution in the lab frame is:

$$f(\vec{v}) = \frac{1}{(2\pi\sigma_v^2)^{3/2}} \exp\left(-\frac{(\vec{v} + \vec{v}_\odot)^2}{2\sigma_v^2}\right)$$

$$\left\{ \begin{array}{l} \vec{v}_\odot : \text{Sun's velocity} \\ \sigma_v = v_0/\sqrt{2} : \text{dispersion} \\ v_0 : \text{circular speed at Solar radius} \end{array} \right. \quad v_0 = 220 \text{ km/s}$$

But the Halo can be:

- ellipsoidal, triaxial (change  $f(v)$ ), (co)-rotating
- anisotropic, ...

The distribution of mass in the halo can be described by the Navarro-Frenk-White profile:

$$\frac{\rho(r)}{\rho_{\text{crit}}} = \frac{\delta_c}{(r/r_s)(1 + r/r_s)^2}$$

NFW 1996 ApJ 462, 563

$r_s$  scale radius

with the critical density

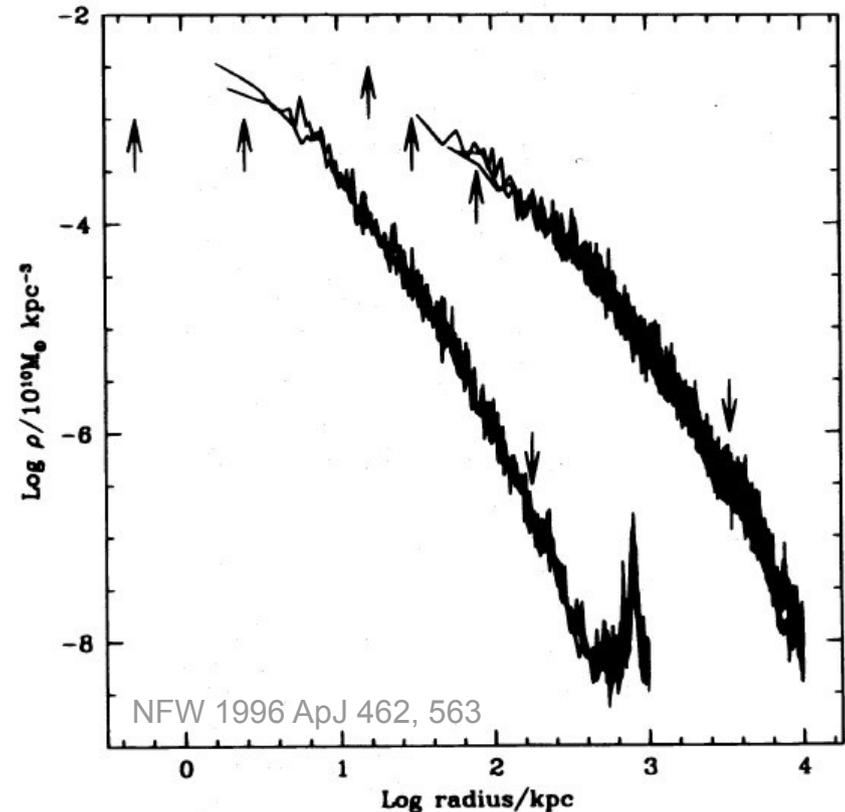
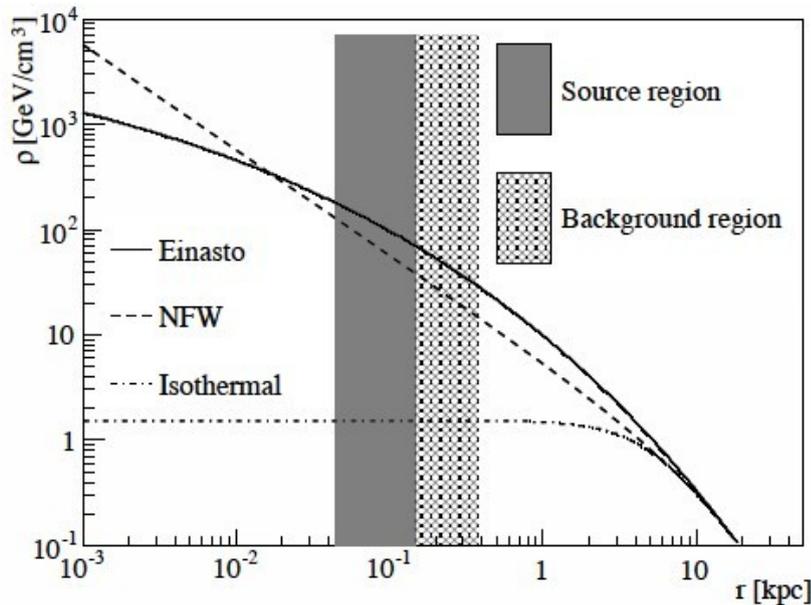
$$\rho_{\text{crit}} = \frac{3H^2}{8\pi G}$$

$\delta_c$  is a dimensionless parameter

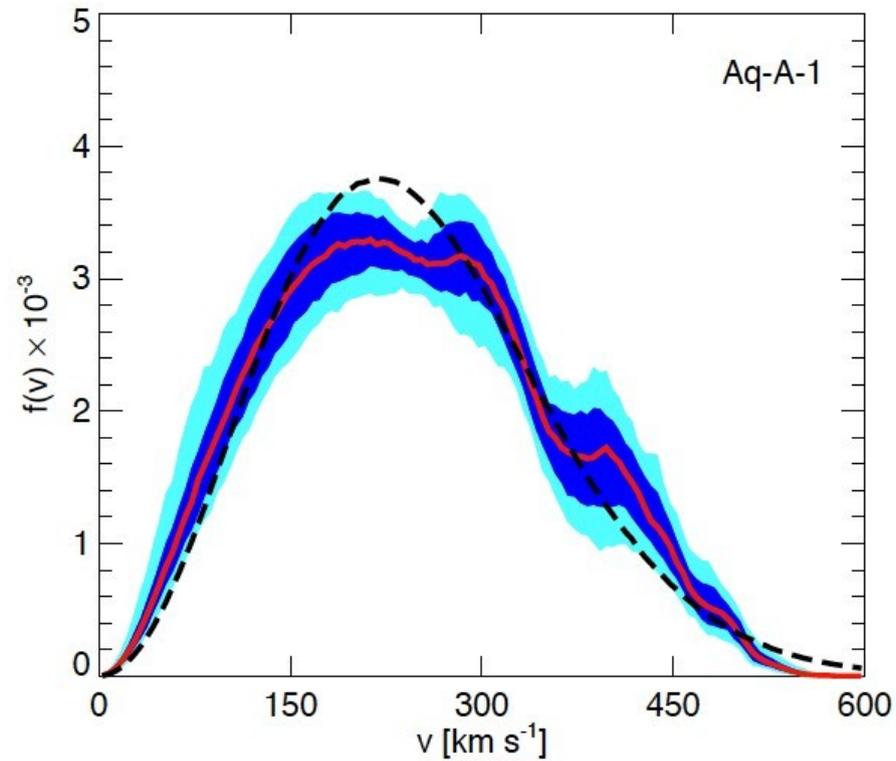
# Expected DM Halo Distribution – Spatial

- Average density may be described by NFW formula
- Lumpiness:
  - ▶ Numerical simulations do not resolve scales tested by direct detection experiments
  - ▶ Baryonic matter dominates over DM within the galactic disk.  
Result: subhalo structures should be smoothed out within the disk.

$$\frac{\rho(r)}{\rho_{\text{crit}}} = \frac{\delta_c}{(r/r_s)(1 + r/r_s)^2}$$

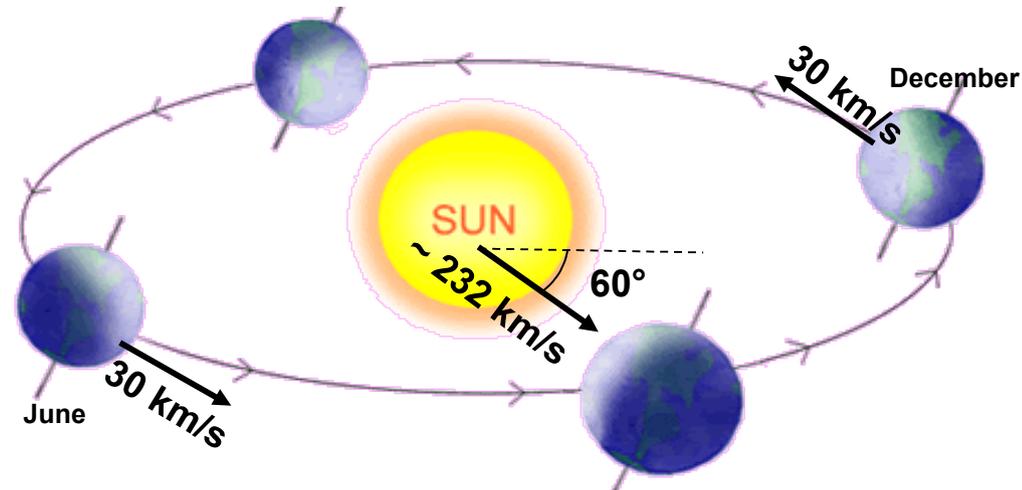


# Expected DM Halo Distribution – Velocity



Vogelsberger et al., 2008 0812.0362

# Annual Modulation



$$\frac{dR}{dE_R} \approx \left( \frac{dR}{dE_R} \right) [1 + \Delta(E_R) \cos \alpha(t)]$$

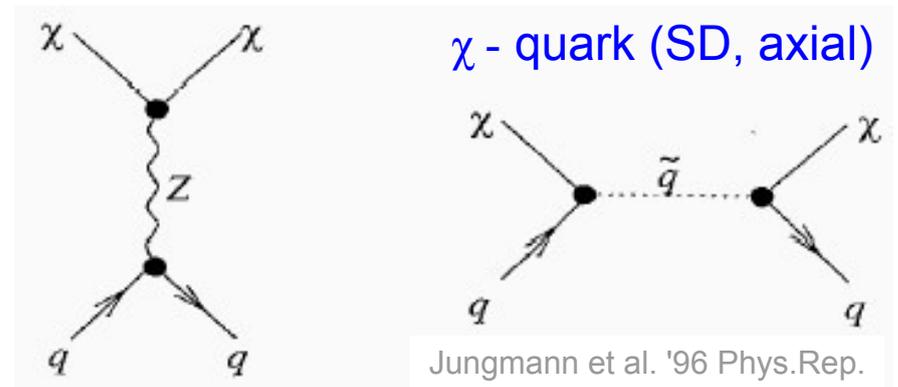
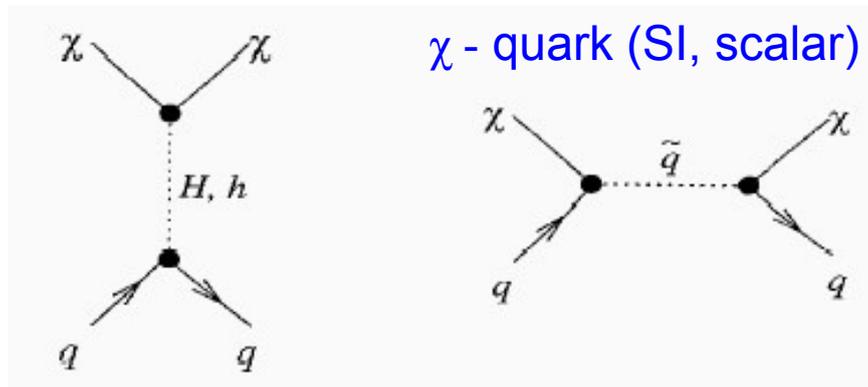
where  $\alpha(t) = 2\pi(t - t_0)/T$ ,  $T = 1$  year and  $t_0 \sim 150$  days.

# Particle Physics Input

# WIMP Scattering Cross Sections

**Example:** SUSY (but direct searches are sensitive to other models as well)

- Compute cross sections  $\chi$  – quark and  $\chi$  – gluon with various SUSY models. Large parameter space, constrained by accelerator and direct search experiments, and cosmology.
  - ▶ **spin-independent:** coupling to mass of nucleus. Coherence  $\Rightarrow \sigma \propto A^2$
  - ▶ **spin-dependent:** coupling of spins of nucleus and neutralino interaction with paired nucleons in the same energy state cancel  $\Rightarrow$  no  $A^2$  enhancement

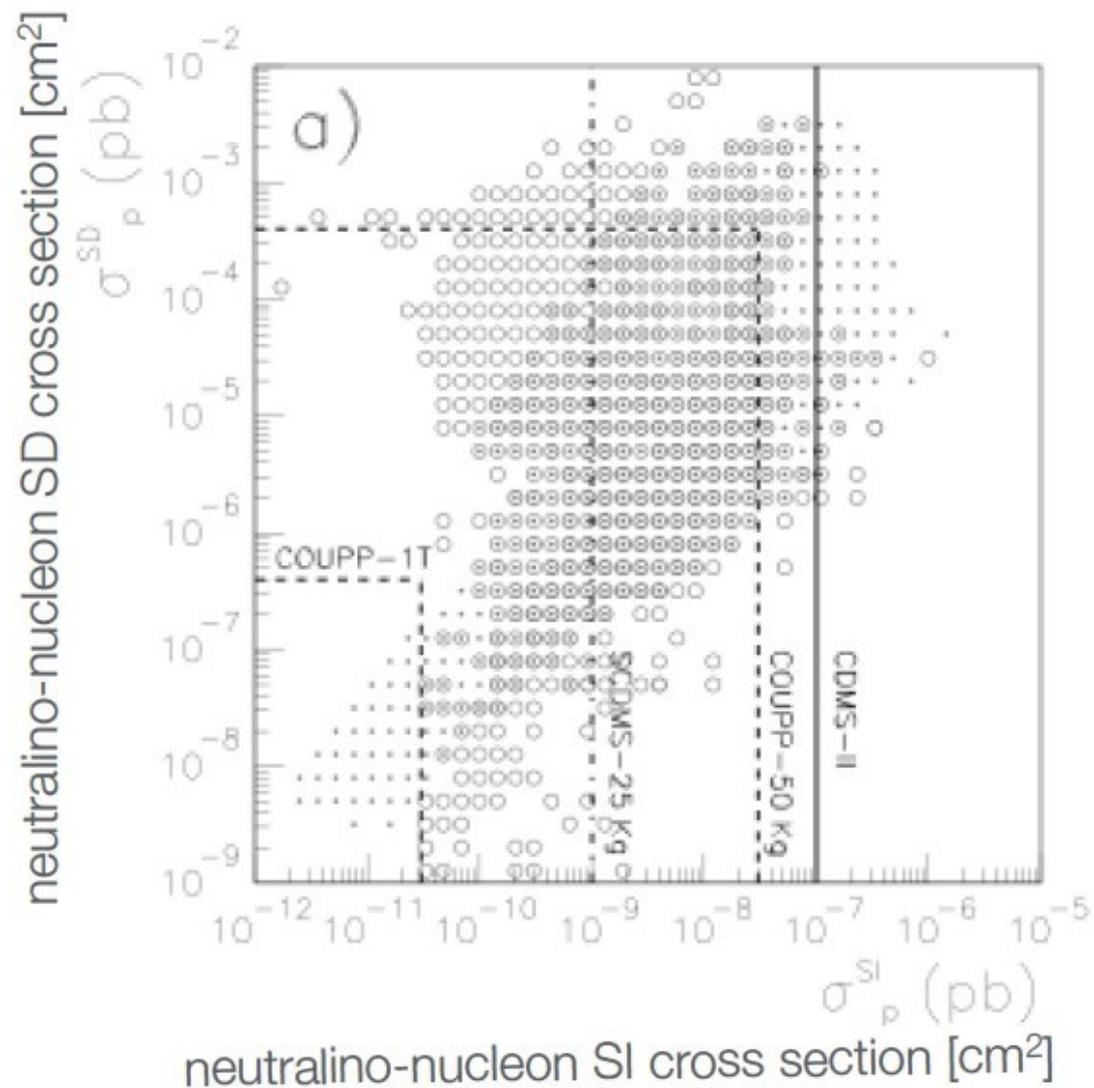


- Distribution of nucleons within nucleus: nuclear form factor.
  - ▶ SI: Large nuclei gain  $\sim A^2$  at small momentum transfer, but lose at higher momentum transfer due to coherence loss.

# Cross Section for WIMP-Nucleon Scattering Spin-Dependent

$$\sigma_{SD} = \frac{32}{\pi} G_F^2 \frac{m_\chi^2 m_N^2}{(m_\chi + m_N)^2} \frac{J_N + 1}{J_N} \left( a_p \langle S_p \rangle + a_n \langle S_n \rangle \right)^2$$

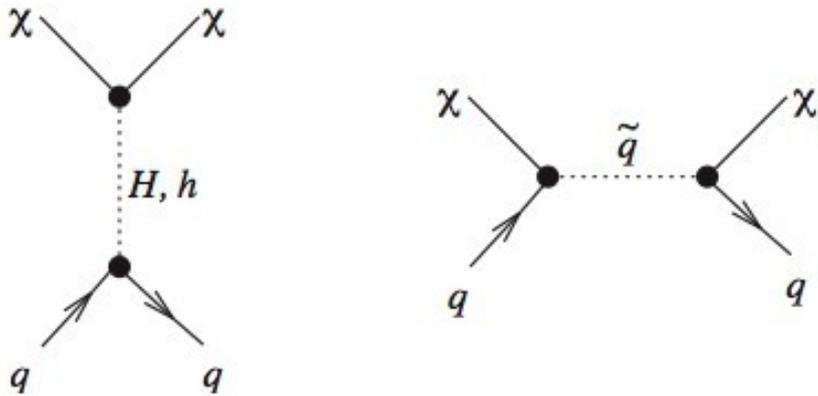
# Spin-dependent vs. Spin-Independent Scattering in the CMSSM



# Nuclear Physics Input

# Cross Section for WIMP-Nucleon Scattering, Nuclear Form Factor

Scalar (spin-independent) scattering

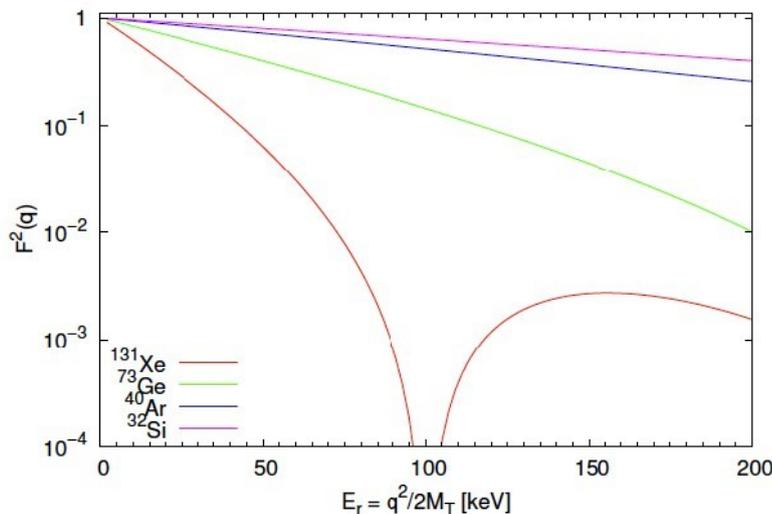


At zero momentum transfer:

$$\sigma = \frac{4 \mu^2}{\pi} [Z f_p + (A - Z) f_n]^2 \quad \text{coherent scattering}$$

reduced mass  $\mu = \frac{m_\chi m_N}{m_\chi + m_N}$   $m_\chi \ll m_N : \mu \approx m_\chi$   
 $m_\chi \gg m_N : \mu \approx m_N$

$f_p, f_n$  : scattering amplitudes protons, neutrons



At finite momentum transfer:

**Form factor:**

Fourier transform of nuclear density

$$F^2(q) = \left| \int \rho(r) \exp \left\{ i \frac{\vec{q} \cdot \vec{r}}{\hbar} \right\} dr \right|^2$$

$$F(E_r) = \left( \frac{3 j_1(qR_1)}{(qR_1)} \right)^2 \exp[-(qs)^2]$$

Momentum transfer:  $q = \sqrt{s m_N E_r}$

$j_1$  : first spherical Bessel function

$$R_1 = \sqrt{R_0^2 - 5 s^2}$$

$$R_0 \approx 1.2 \text{ fm } A^{1/3}$$

$$s \approx 1 \text{ fm}$$

Helms form factor based on:

$$\rho(r) = \int_{\text{volume}} \rho_0(\mathbf{r}') \rho_1(\mathbf{r} - \mathbf{r}') d^3 x'$$

$$\rho_0(r) = \begin{cases} \frac{3}{4\pi r_n^3} & r < r_n \\ 0 & r > r_n, \end{cases}$$

$$\rho_1(r) = \frac{1}{(2\pi s^2)^{3/2}} e^{-r^2/2s^2}.$$

# Recoil Spectrum

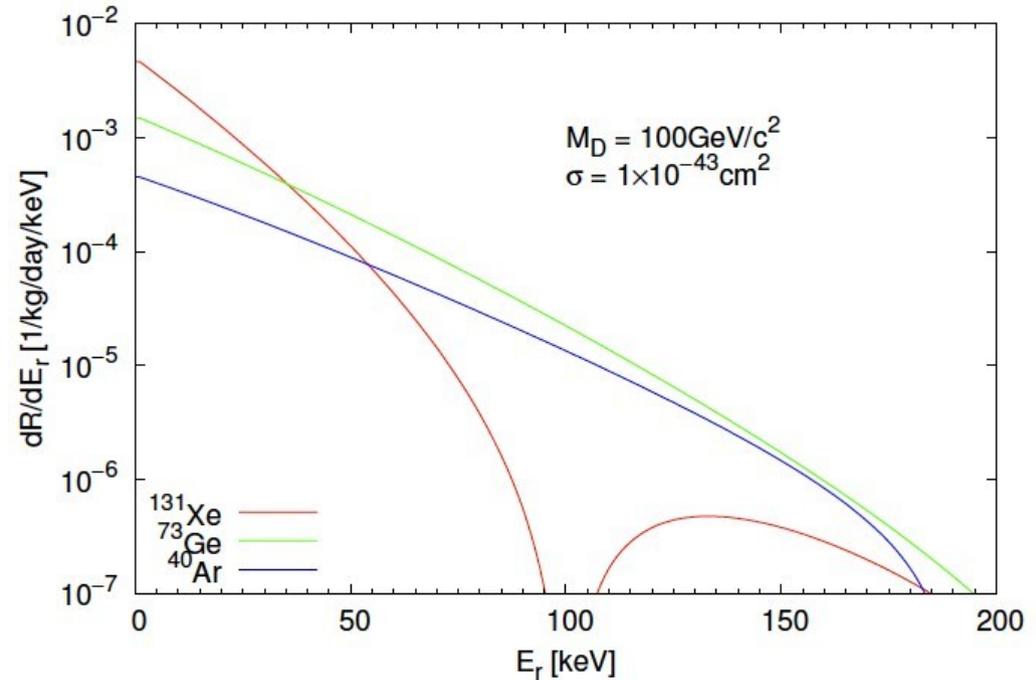
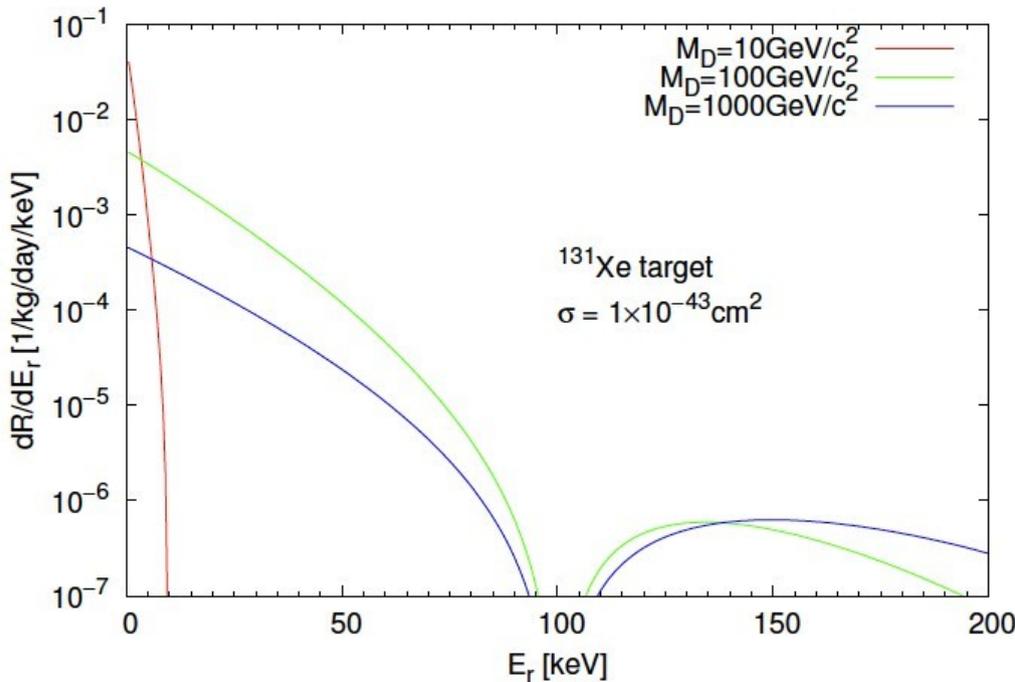
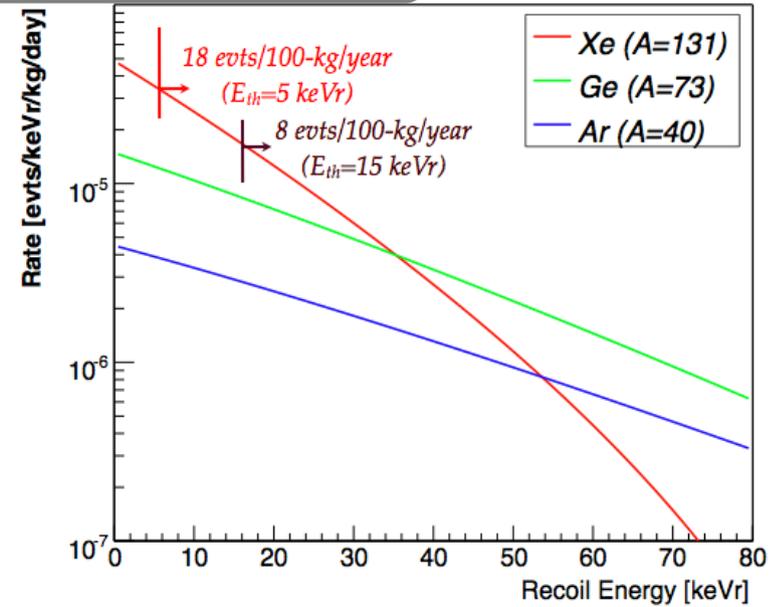
$$\frac{dR}{dE} = \frac{\rho_\chi \sigma_s}{2 m_\chi \mu^2} |F(E)^2| \int_{v_{min}}^{v_{esc}} f(\mathbf{v}, t) \frac{(\mathbf{v}, t)}{v} d^3 v$$

$$f(\mathbf{v}, t) \propto \exp\left(-\frac{(\mathbf{v} + \mathbf{v}_E(t))^2}{2 \sigma_v^2}\right)$$

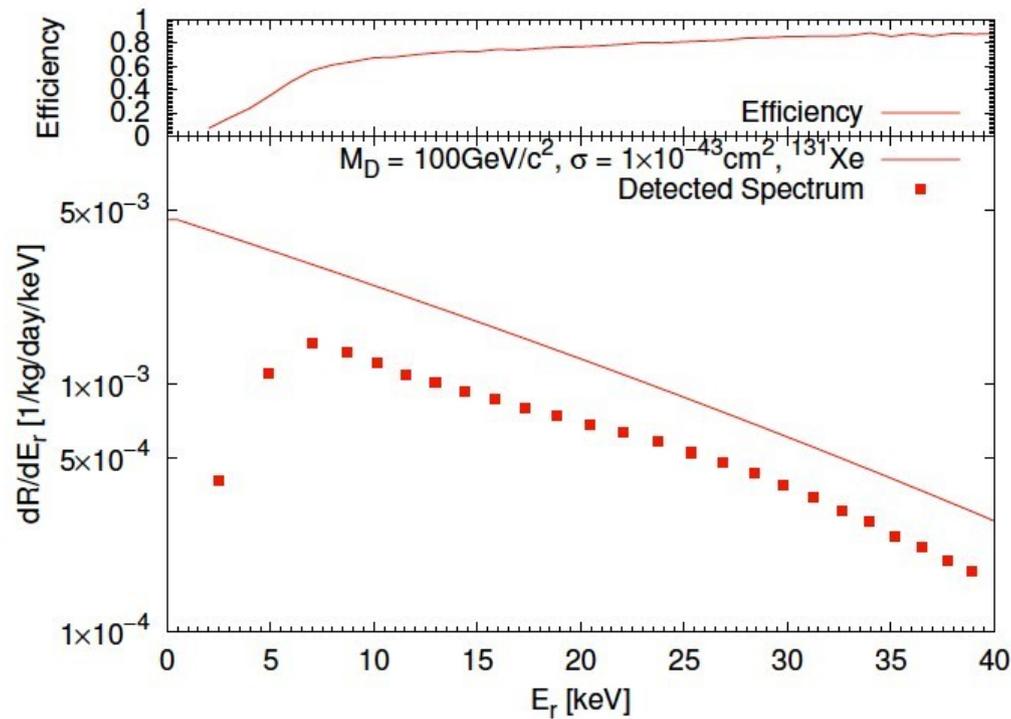
- For standard halo model approximate spectral form:

$$\frac{dR}{dE_R} \approx \left(\frac{dR}{dE_R}\right)_0 F^2(E_R) \exp\left(-\frac{E_R}{E_c}\right)$$

- Including form factors:



# Detector Effects: Smearing of the Threshold, Efficiencies



# Backgrounds in Direct DM Search

Cross-sections are very small:  $<10^{-43}$  cm<sup>2</sup> or  $10^{-7}$  pb (spin-independent)

Without background, sensitivity  $\propto$  (mass  $\times$  exposure time)<sup>-1</sup>

With background subtraction  $\propto$  (M t)<sup>-1/2</sup>  
until limited by systematics.

Backgrounds:

**Gamma-rays & beta decays:**

~100 events/kg/day

Need very good  $\beta$  and  $\gamma$  background discrimination.

Shielding: low-activity lead, water, noble liquids (active), liquid N<sub>2</sub>, ...

**Neutrons from ( $\alpha$ , n) and spontaneous fission (concrete, rock, etc.):**

~ 1 event/kg/day (LNGS)

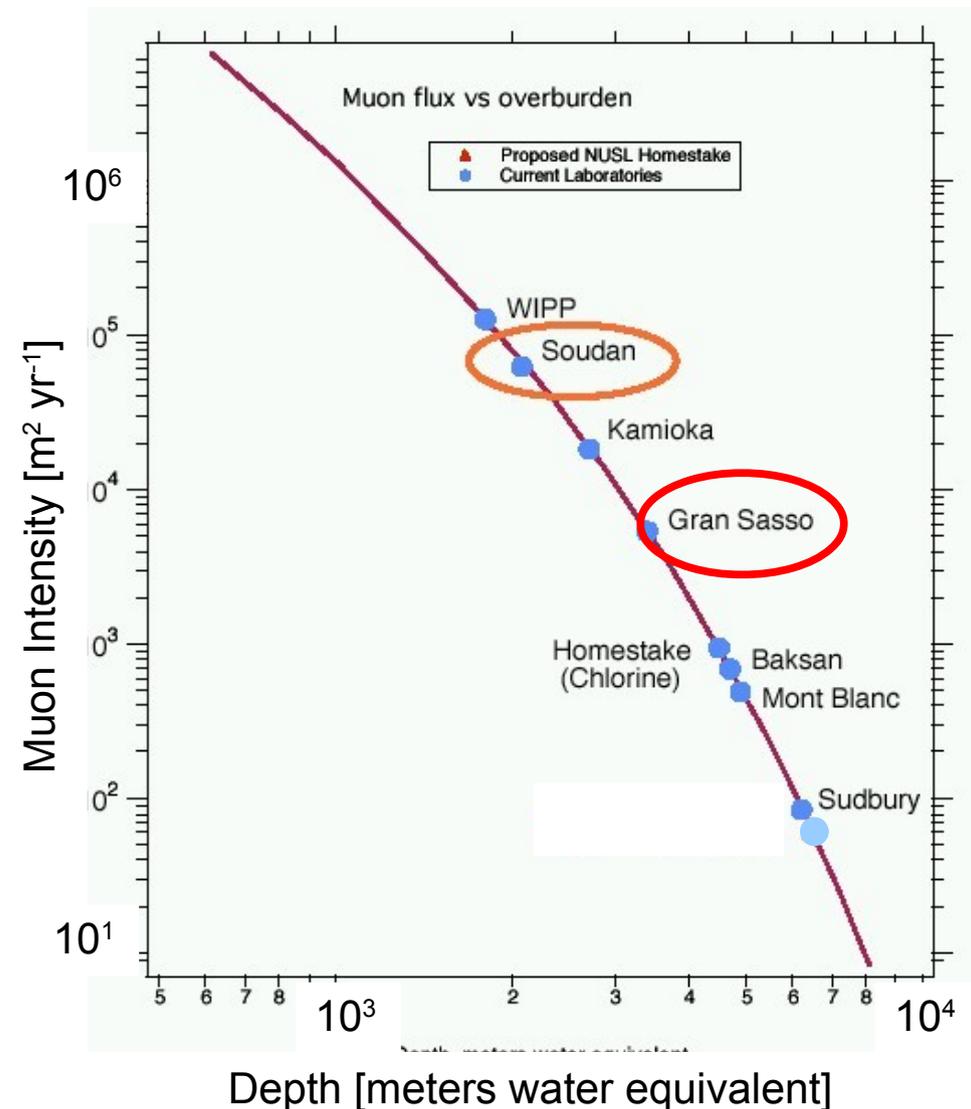
Neutron moderator (polyethylene, paraffin, ...)

**Neutrons from CR muons:**

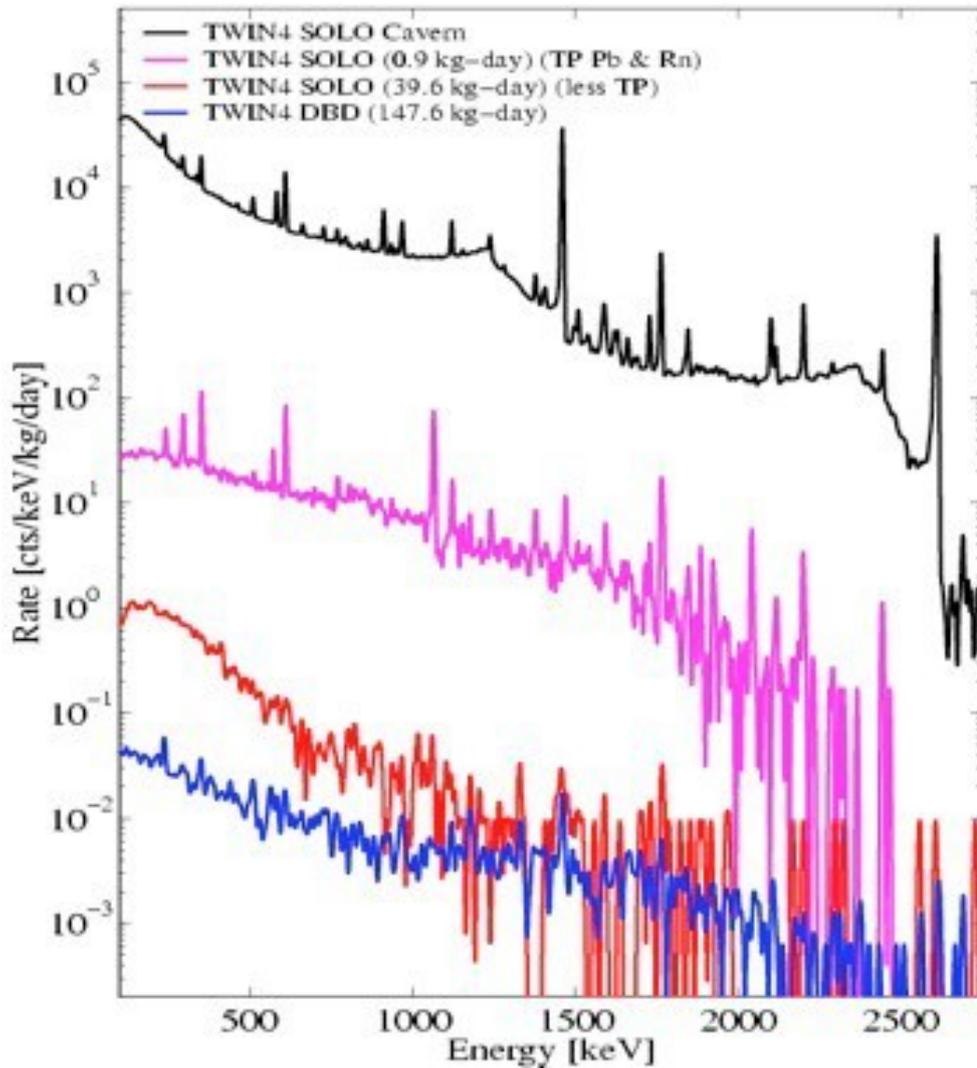
Rate depending on depth.

$\mu$ -veto, n-veto, shielding

**$\alpha$  decays from Rn daughters, ...**



# Backgrounds in Direct DM Search

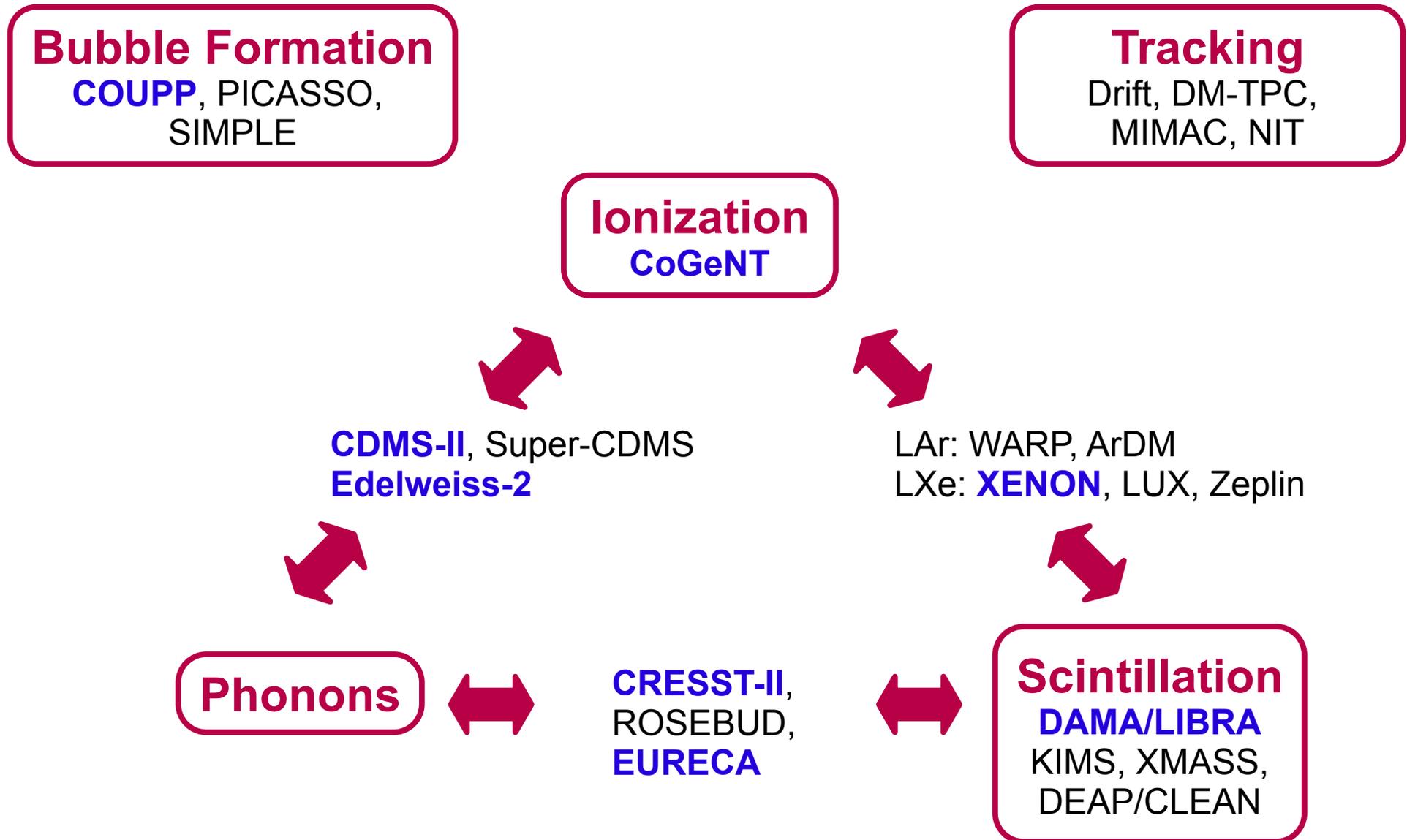


- Ge spectrum unshielded underground

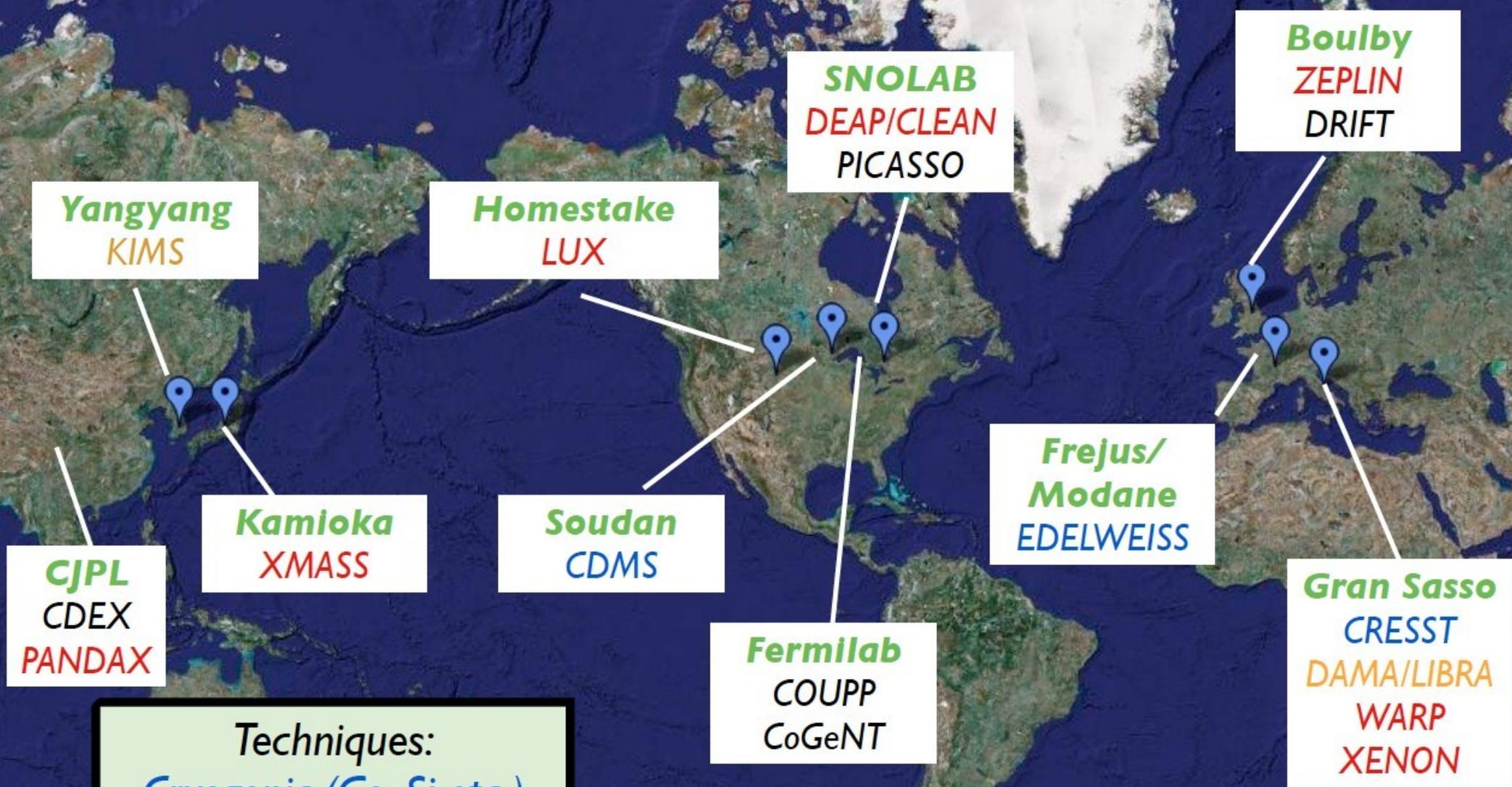
- Ge spectrum underground with Pb shield and purge for Rn

# DM Detector Overview

## Detection Principles



# Looking for Dark Matter at Underground Labs



**Techniques:**  
Cryogenic (Ge, Si etc.)  
Solid Scintillator (NaI, CsI)  
Noble Liquids (LXe, LAr)